

# Quantum Algorithms Tutorial

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# Post-quantum cryptography

- ▶ Quantum computers can break public-key cryptography that is based on assuming hardness of factoring, discrete logs, and a few other problems
- ▶ Post-quantum cryptography tries to design classical crypto schemes that cannot be broken efficiently by quantum algorithms
- ▶ Classical codemakers vs quantum codebreakers
- ▶ This tutorial:

Get to know your enemy!

# Quantum bits

- ▶ Richard Feynman,  
David Deutsch  
in early 1980s:



Harness quantum effects for useful computations!

- ▶ Classical bit is 0 or 1; **quantum bit** is superposition of 0 and 1  
For example, can use an electron as qubit,  
with 0 = “spin up”    and    1 = “spin down”
- ▶ 2 qubits is superposition of **4** basis states (00,01,10,11)  
3 qubits is superposition of **8** basis states (000,001, ... )  
...  
1000 qubits: superposition of  **$2^{1000}$**  states
- ▶ Massive space for computation! Easier said than done...

## A bit of math: states

► 1-qubit basis states:  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

► Qubit: superposition  $\alpha_0|0\rangle + \alpha_1|1\rangle = \begin{pmatrix} \alpha_0 \\ \alpha_1 \end{pmatrix} \in \mathbb{C}^2$

►

2-qubit basis state:  $|10\rangle = |1\rangle \otimes |0\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$

►  $n$ -qubit state:  $|\psi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle \in \mathbb{C}^{2^n}$

► Axiom: measuring state  $|\psi\rangle$  gives  $|x\rangle$  with probability  $|\alpha_x|^2$

► Hence  $\sum_{x \in \{0,1\}^n} |\alpha_x|^2 = 1$ , so  $|\psi\rangle$  is a vector of length 1

# A bit of math: operations

- ▶ Quantum operation maps quantum states to quantum states and is *linear*  $\implies$  corresponds to **unitary** matrix

- ▶ Example 1-qubit gates:

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 0 \\ 0 & e^{\pi i/4} \end{pmatrix}$$

- ▶ More quantum: **Hadamard** gate  $= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle), \quad H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

$$\text{But } H \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) = \frac{1}{\sqrt{2}} H|0\rangle + \frac{1}{\sqrt{2}} H|1\rangle = |0\rangle$$

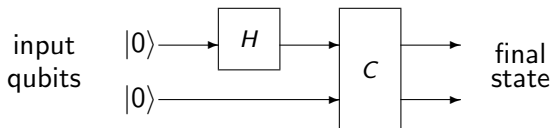
**Interference!**

- ▶ **Controlled-NOT** gate on 2 qubits:  $|a, b\rangle \mapsto |a, a \oplus b\rangle$

# Quantum circuits

- ▶ A classical Boolean circuit consists of AND, OR, and NOT gates on an  $n$ -bit register
- ▶ A **quantum circuit** consists of unitary **quantum gates** on an  $n$ -qubit register (allowing  $H$ ,  $T$ , and CNOT gates suffices)

Example:



$$|00\rangle \xrightarrow{H \otimes I} \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle) \xrightarrow{\text{CNOT}} \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

This circuit creates an EPR-pair: **entanglement!**

## Recap: From classical to quantum computation

- ▶ bits  $\longrightarrow$  qubits
- ▶ AND/OR/NOT gates  $\longrightarrow$  unitary quantum gates
- ▶ classical circuit  $\longrightarrow$  quantum circuit
- ▶ reading the output bit  $\longrightarrow$  measuring final state

# Quantum mechanical computers

1. Start with all qubits in easily-preparable state (e.g. all  $|0\rangle$ )
2. Run a circuit that produces the right kind of interference: computational paths leading to correct output should interfere constructively, others should interfere destructively
3. Measurement of final state gives classical output

Two important questions:

1. Can we build such a computer?
2. What can it do?

This tutorial: 2nd question, focus on quantum algorithms



# Quantum parallelism

- ▶ Suppose classical algorithm computes  $f : \{0,1\}^n \rightarrow \{0,1\}^m$
- ▶ Convert this to quantum circuit  $U : |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle$
- ▶ We can now compute  $f$  “on all inputs **simultaneously**”!

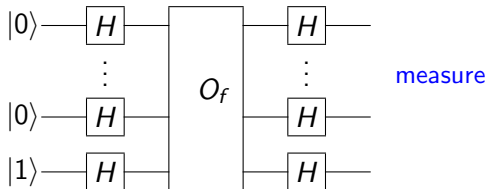
$$U \left( \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle|0\rangle \right) = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle|f(x)\rangle$$

- ▶ This contains all  $2^n$  function values!
- ▶ But observing gives only one random  $|x\rangle|f(x)\rangle$   
All other information will be lost
- ▶ **More tricks needed for successful quantum computation**  
**Interference!**

# Deutsch-Jozsa problem

- ▶ Given: function  $f : \{0, 1\}^n \rightarrow \{0, 1\}$  ( $2^n$  bits) , s.t.  
(1)  $f(x) = 0$  for all  $x$  (**constant**),  
or  
(2)  $f(x) = 0$  for  $\frac{1}{2} \cdot 2^n$  of the  $x$ 's (**balanced**)
- ▶ Question: is  $f$  constant or balanced?
- ▶ **Classically**: need at least  $\frac{1}{2} \cdot 2^n + 1$  steps ("queries" to  $f$ )
- ▶ **Quantumly**:  $O(n)$  gates suffice, and only 1 query
- ▶ Query: application of unitary  $O_f : |x, 0\rangle \mapsto |x, f(x)\rangle$
- ▶ More generally:  $O_f : |x, b\rangle \mapsto |x, b \oplus f(x)\rangle$  ( $b \in \{0, 1\}$ )
- ▶ NB using  $|-\rangle = H|1\rangle$ , we can get queried bit as a  $\pm$ -phase:  
 $O_f|x\rangle|-\rangle = (-1)^{f(x)}|x\rangle|-\rangle$

# Deutsch-Jozsa algorithm



- ▶ Starting state:  $\underbrace{|0 \dots 0\rangle}_n |1\rangle$

- ▶ After first **Hadamards**:  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |-\rangle$

- ▶ **Make one query**:  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} |x\rangle |-\rangle$

- ▶ Forget about the last qubit  $|-\rangle$

## Deutsch-Jozsa (continued)

- ▶ After second Hadamard:

$$\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} (-1)^{x \cdot y} |y\rangle$$

- ▶  $\alpha_{0\dots 0} = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} = \begin{cases} 1 & \text{if constant} \\ 0 & \text{if balanced} \end{cases}$

- ▶ Measurement gives right answer with certainty
- ▶ Big quantum-classical separation:  $O(n)$  vs  $\Omega(2^n)$  steps
- ▶ But the problem is efficiently solvable by bounded-error classical algorithm (just query  $f$  at a few random  $x$ )

# The meat of this tutorial: 4 quantum algorithms

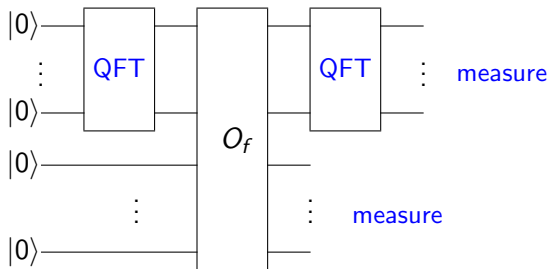
1. Shor's factoring algorithm
2. Grover's search algorithm
3. Ambainis's collision-finding algorithm
4. HHL algorithm for linear systems

# Factoring

- ▶ Given  $N = p \cdot q$ , compute the prime factors  $p$  and  $q$
- ▶ Fundamental **mathematical** problem since Antiquity
- ▶ Fundamental **computational** problem on  $\log N$  bits  
 $15 = 3 \times 5$   
 $12140041 = 3413 \times 3557$
- ▶ Best known classical algorithms use time  $2^{(\log N)^\alpha}$ , where  $\alpha = 1/2$  or  $1/3$
- ▶ Its **assumed** computational hardness is basis of **public-key cryptography** (RSA)
- ▶ A quantum computer can **break** this, using **Shor's efficient quantum factoring algorithm**!

# Overview of Shor's algorithm

- ▶ Classical reduction: choose random  $x \in \{2, \dots, N - 1\}$ .  
It suffices to find **period**  $r$  of  $f(a) = x^a \bmod N$
- ▶ Shor uses the **quantum Fourier transform** for period-finding



- ▶ Overall complexity: roughly  $(\log N)^2$  elementary gates

## Reduction to period-finding

- ▶ Pick a random integer  $x \in \{2, \dots, N-1\}$ , s.t.  $\gcd(x, N)=1$
- ▶ The sequence  $x^0, x^1, x^2, x^3, \dots \bmod N$  cycles:  
has an unknown **period**  $r$  (min  $r > 0$  s.t.  $x^r \equiv 1 \bmod N$ )
- ▶ With probability  $\geq 1/4$  (over the choice of  $x$ ):  
 $r$  is even and  $x^{r/2} \pm 1 \not\equiv 0 \bmod N$
- ▶ Then:  
$$\begin{aligned}x^r &= (x^{r/2})^2 \equiv 1 \bmod N \iff \\(x^{r/2} + 1)(x^{r/2} - 1) &\equiv 0 \bmod N \iff \\(x^{r/2} + 1)(x^{r/2} - 1) &= kN \text{ for some } k\end{aligned}$$
- ▶  $x^{r/2} + 1$  and  $x^{r/2} - 1$  each share a factor with  $N$
- ▶ This factor of  $N$  can be extracted using gcd-algorithm



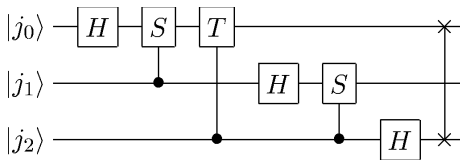
# Quantum Fourier transform

- **Fourier basis** (dimension  $q$ ):  $|\chi_j\rangle = \frac{1}{\sqrt{q}} \sum_{k=0}^{q-1} e^{\frac{2\pi i j k}{q}} |k\rangle$

Such a state is unentangled  $|\chi_{j_0 j_1 j_2}\rangle =$

$$\frac{1}{\sqrt{8}} (|0\rangle + e^{2\pi i 0 \cdot j_2} |1\rangle) \otimes (|0\rangle + e^{2\pi i 0 \cdot j_1 j_2} |1\rangle) \otimes (|0\rangle + e^{2\pi i 0 \cdot j_0 j_1 j_2} |1\rangle)$$

- Quantum Fourier Transform:  $|j\rangle \mapsto |\chi_j\rangle$
- If  $q = 2^\ell$ , then can implement this with  $O(\ell^2)$  gates.



- For Shor: choose  $q = 2^\ell$  in  $(N^2, 2N^2]$

## Easy case for the analysis: $r|q$

1. Apply QFT to 1st register of  $\underbrace{|0 \dots 0\rangle}_{\ell \text{ qubits}} \underbrace{|0 \dots 0\rangle}_{\lceil \log N \text{ qubits} \rceil} :$

$$\frac{1}{\sqrt{q}} \sum_{a=0}^{q-1} |a\rangle |0\rangle$$

2. Compute  $f(a) = x^a \bmod N$  (by repeated squaring)

$$\frac{1}{\sqrt{q}} \sum_{a=0}^{q-1} |a\rangle |x^a \bmod N\rangle$$

3. Observing 2nd register gives  $|x^s \bmod N\rangle$  (random  $s < r$ )  
1st register collapses to superposition of

$$|s\rangle, |r+s\rangle, |2r+s\rangle, \dots, |q-r+s\rangle$$

## Easy case: $r|q$ (continued)

Recall: 1st register is in superposition  $\sum_{j=0}^{q/r-1} |jr + s\rangle$

4. Apply QFT once more:

$$\sum_{j=0}^{q/r-1} \sum_{b=0}^{q-1} e^{2\pi i \frac{(jr+s)b}{q}} |b\rangle = \sum_{b=0}^{q-1} e^{2\pi i \frac{sb}{q}} \underbrace{\left( \sum_{j=0}^{q/r-1} \left( e^{2\pi i \frac{rb}{q}} \right)^j \right)}_{\text{geometric sum}} |b\rangle$$

Sum  $\neq 0$  iff  $e^{2\pi i \frac{rb}{q}} = 1$  iff  $\frac{rb}{q}$  is an integer

Only the  $b$  that are multiples of  $\frac{q}{r}$  have non-zero amplitude!

## Easy case: $r|q$ (continued)

5. Observe 1st register: random multiple  $b = c \frac{q}{r}$ ,  $c \in [0, r)$ :

$$\frac{b}{q} = \frac{c}{r}$$

- ▶  $b$  and  $q$  are known;  $c$  and  $r$  are unknown
- ▶  $c$  and  $r$  are coprime with probability  $\geq 1/\log \log r$
- ▶ Then: we find  $r$  by writing  $\frac{b}{q}$  in lowest terms
- ▶ Since we can find  $r$ , we can find prime factors of  $N$  !

Hard case ( $r \nmid q$ ) still works approximately: measurement gives  $b$  s.t.  $\frac{b}{q} \approx \frac{c}{r}$ ; we can find  $r$  with some extra number theory

# Summary for Shor's algorithm

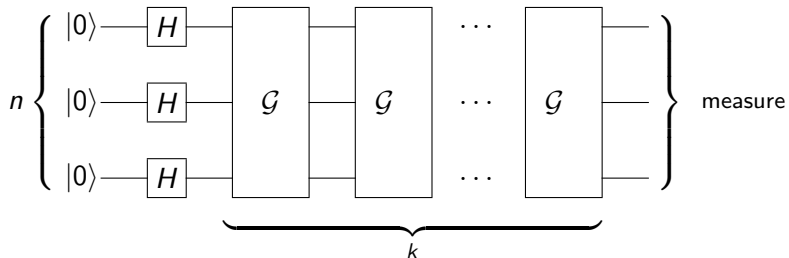
- ▶ Reduce factoring to finding the **period  $r$**  of modular exponentiation function  $f(a) = x^a \bmod N$
- ▶ Use **quantum Fourier transform** to find a multiple of  $q/r$ , repeat a few times to find  $r$
- ▶ Overall complexity:
  - ▶ QFT takes  $O(\log q)^2 = O(\log N)^2$  elementary gates
  - ▶ Modular exponentiation:  $\approx (\log N)^2 \log \log N$  gates; classical computation by repeated squaring (use Schönhage-Strassen algo for fast multiplication)
  - ▶ Everything repeated  $O(\log \log N)$  times
  - ▶ Classical postprocessing takes  $O(\log N)^2$  gates
- ▶ Roughly  **$(\log N)^2$  elementary gates in total**

# The search problem

- ▶ We want to search for some good item in an unordered  $N$ -element search space
- ▶ Model this as function  $f : \{0, 1\}^n \rightarrow \{0, 1\}$  ( $N = 2^n$ )  
 $f(x) = 1$  if  $x$  is a solution
- ▶ We can query  $f$ :  
 $O_f : |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle$   
or  
 $O_f : |x\rangle \mapsto (-1)^{f(x)}|x\rangle$
- ▶ Goal: find a solution
- ▶ Classically this takes  $O(N)$  steps (queries to  $f$ )
- ▶ Grover's algorithm does it in  $O(\sqrt{N})$  steps

# Grover's algorithm

- ▶ Apply Grover iteration  $\mathcal{G}$   $k$  times on uniform starting state



- ▶ Idea: each iteration moves amplitude towards solutions

# The good state and the bad state

- ▶ Suppose there are  $t$  solutions
- ▶ Define “good” state and “bad” state:

$$|G\rangle = \frac{1}{\sqrt{t}} \sum_{x:f(x)=1} |x\rangle \quad |B\rangle = \frac{1}{\sqrt{N-t}} \sum_{x:f(x)=0} |x\rangle$$

- ▶ Initial uniform state is  $|U\rangle = \sin(\theta)|G\rangle + \cos(\theta)|B\rangle$   
for  $\theta = \arcsin(\sqrt{t/N})$
- ▶ All intermediate states will be in  $\text{span}\{|G\rangle, |B\rangle\}$
- ▶ Grover iteration is a rotation over angle  $2\theta$   
so after  $k$  iterations the state is

$$\sin((2k+1)\theta)|G\rangle + \cos((2k+1)\theta)|B\rangle$$



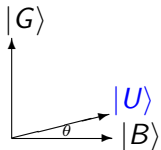
# One Grover iteration: rotation by $2\theta$

$\mathcal{G} = H^{\otimes n} R H^{\otimes n} \cdot O_f$ , where  $R$  reflects through  $|0^n\rangle$

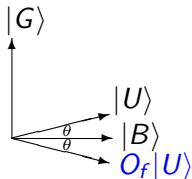
This  $\mathcal{G}$  is the **product of two reflections**:

1.  $O_f$  reflects through  $|B\rangle$
2.  $H^{\otimes n} R H^{\otimes n}$  reflects through  $|U\rangle$

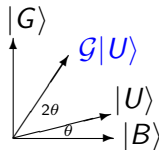
Starting state:



Reflect through  $|B\rangle$ :



Reflect through  $|U\rangle$ :



## How many iterations do we need?

- ▶ Success probability after  $k$  iterations:

$$\sin^2((2k+1)\theta), \text{ with } \theta = \arcsin(\sqrt{t/N}) \approx \sqrt{t/N}$$

- ▶ If  $k = \frac{\pi}{4\theta} - \frac{1}{2}$ , then success probability is  $\sin^2(\pi/2) = 1$
- ▶ Example:  $t = N/4$  solutions  $\Rightarrow k = 1$
- ▶ In general, round  $k$  to nearest integer (incurs small error)
- ▶ Query complexity is  $k \approx \frac{\pi}{4} \sqrt{N/t}$   
This is optimal for a quantum algorithm!
- ▶ Gate complexity is  $O(\sqrt{N/t} \log N)$

# Summary for Grover's algorithm

- ▶ Quantum computers can search any  $N$ -element space with  $t = \varepsilon N$  solutions, in  $O(\sqrt{N/t}) = O(1/\sqrt{\varepsilon})$  iterations
  1. Set up uniform starting state  $|U\rangle$
  2. Repeat the following  $O(1/\sqrt{\varepsilon})$  times:
    - 2.1 Reflect through  $|B\rangle$  (costs 1 query)
    - 2.2 Reflect through  $|U\rangle$  (costs  $O(\log N)$  gates)
  3. Measure final state to obtain an index  $i$
- ▶ If we don't know  $\varepsilon = t/N$ , we can try different guesses, still find a solution with expected number of  $O(1/\sqrt{\varepsilon})$  iterations
- ▶ The algorithm has a small error probability, but can be modified to error 0 *if* we know  $t$  exactly

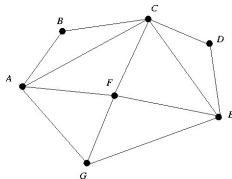
# Application: Speed up NP problems

- ▶ Given a propositional formula  $f(x_1, \dots, x_n)$   
Computable in time  $\text{poly}(n)$

Question: is  $f$  satisfiable?

- ▶ This is a typical NP-complete problem
- ▶ Search space of  $N = 2^n$  possibilities
- ▶ Classically: exhaustive search is the best we know.  
This takes about  $N$  steps
- ▶ Quantumly: Grover finds a satisfying assignment in  $\sqrt{N} \cdot \text{poly}(n)$  steps
- ▶ Because Grover is optimal, we believe that NP-hard problems cannot be efficiently computed by quantum algorithms

# Classical random walks



- ▶ Explore a graph by moving to random neighbor in each step
- ▶ If  $G$  is  $d$ -regular and connected: normalized adjacency matrix has “spectral gap”  $\delta \in (0, 1)$ . Starting from any vertex,  $O(1/\delta)$  random walk steps produce uniform distribution
- ▶ Suppose an  $\varepsilon$ -fraction of the vertices are “marked” and we want to find such a marked vertex. Simple classical algorithm:
  1. Start at random vertex  $v$  (setup cost  $\mathbf{S}$ )
  2. Do the following  $O(1/\varepsilon)$  times:
    - 2.1 Check if  $v$  is marked (checking cost  $\mathbf{C}$ )
    - 2.2 Rerandomize  $v$  by  $O(1/\delta)$  RW steps (step cost  $\mathbf{U}$ )

This finds marked item w.h.p. Cost is  $\mathbf{S} + \frac{1}{\varepsilon} \left( \mathbf{C} + \frac{1}{\delta} \mathbf{U} \right)$

# Quantum walks

- ▶ **Quantum walk**: walk in superposition over vertices (edges)
- ▶ Analogy with Grover's algorithm:
  - $|G\rangle$  = uniform superposition over edges with marked endpoint
  - $|B\rangle$  = uniform superposition over all other edges
  - $|U\rangle = \sin(\theta)|G\rangle + \cos(\theta)|B\rangle$ ,  $\theta = \arcsin(1/\sqrt{\epsilon})$
- 1. Setup starting state  $|U\rangle$  (setup cost **S**)
- 2. Repeat the following  $O(1/\sqrt{\epsilon})$  times:
  - 2.1 Reflect through  $|B\rangle$  (checking cost **C**)
  - 2.2 Reflect through  $|U\rangle$   
(can be implemented using  $1/\sqrt{\delta}$  QW steps, each at cost **U**)
- 3. Measure and check that resulting vertex is marked.

Correctness analogous to Grover. Cost is  $\mathbf{S} + \frac{1}{\sqrt{\epsilon}} \left( \mathbf{C} + \frac{1}{\sqrt{\delta}} \mathbf{U} \right)$

## Example: Ambainis's algorithm ('03)

Suppose we want to find a collision in  $h : [n] \rightarrow \mathbb{N}$

- ▶  $G = \text{Johnson graph}$ : the vertices are the sets  $R \subseteq [n]$  of size  $r$ .  
Edge between sets  $R$  and  $R'$  if they differ in 1 element
- ▶ Fraction of vertices of  $G$  that contain collision:  $\epsilon \geq (r/n)^2$
- ▶ Known: spectral gap is  $\delta \approx 1/r$
- ▶ With each vertex  $R$ , algorithm records  $h(R)$ ;  
setup cost  $\mathbf{S} = r$ ; checking cost  $\mathbf{C} = 0$ ; update cost  $\mathbf{U} = O(1)$
- ▶ Total cost:  $\mathbf{S} + \frac{1}{\sqrt{\epsilon}} \left( \mathbf{C} + \frac{1}{\sqrt{\delta}} \mathbf{U} \right) \stackrel{r=n^{2/3}}{=} O(n^{2/3})$
- ▶ Classically:  $\Theta(n)$   $f$ -evaluations needed

If  $h$  is 2-to-1: run on random set of  $\sqrt{n}$  inputs (whp 1 collision) to get complexity  $O(n^{1/3})$

Classically:  $\Theta(\sqrt{n})$   $f$ -evaluations, by birthday paradox

# HHL algorithm for “solving” large linear systems

- ▶ Solving large linear systems  $Ax = b$  is one of the most important problems in science and engineering.

Goal: given matrix  $A$  and vector  $b$ , find vector  $x$

- ▶ Harrow-Hassidim-Lloyd'09: “solves” this problem **exponentially faster** by preparing state  $|x\rangle$  **IF** system is well-behaved:

Assumptions

- (1) state  $|b\rangle$  easy to prepare;
- (2)  $A$  is well-conditioned:  $\lambda_{\max}/\lambda_{\min}$  not too big;
- (3) unitary operation  $e^{iA}$  is easy to apply (sparseness suffices)



# How does the Harrow-Hassidim-Lloyd algorithm work?

- ▶ Input: Hermitian matrix  $A \in \mathbb{R}^{N \times N}$  and vector  $b \in \mathbb{R}^N$   
Goal: approximately prepare  $|x\rangle$ , where  $Ax = b$
- ▶ Let  $v_1, \dots, v_N, \lambda_1, \dots, \lambda_N$  be eigenvectors, eigenvalues of  $A$
- ▶ HHL algorithm:
  1. Prepare quantum state  $|b\rangle = \sum_{i=1}^N \beta_i |v_i\rangle$   
NB: applying  $A^{-1}$  corresponds to multiplying with  $\lambda_i^{-1}$
  2. Use eigenvalue estimation:  $\sum_{i=1}^N \beta_i |v_i\rangle |\lambda_i\rangle$
  3. Make new qubit  $\sum_{i=1}^N \beta_i |v_i\rangle |\lambda_i\rangle \left( \lambda_i^{-1} |0\rangle + \sqrt{1 - \lambda_i^{-2}} |1\rangle \right)$
  4. Uncompute  $|\lambda_i\rangle$  by inverting eigenvalue estimation
  5. Amplify the  $|0\rangle$ -part to end with  $\sum_{i=1}^N \beta_i \lambda_i^{-1} |v_i\rangle = |x\rangle$

# What else can a quantum computer do?

- ▶ **Similar to Shor:** discrete logarithm, solve Pell's equation, compute properties of number fields, ...
- ▶ **Similar to Grover:** maximum-finding, approximate counting, shortest paths in graphs, minimum spanning trees, ...
- ▶ **Similar to quantum walks:** finding small subgraphs, matrix-product verification, junta-testing, backtracking, ...
- ▶ **Similar to HHL:** quantum machine learning, principal component analysis, recommendation systems, ...
- ▶ Efficiently **simulating quantum-mechanical systems.**  
Could be very important for drug design, material sciences. ...

# What quantum algorithms *cannot* do

- ▶ You can simulate every quantum algorithm with an exponentially slower classical computer

This implies that the set of *computable* problems doesn't change: Church-Turing thesis remains intact

- ▶ For many problems we can show that quantum computers give no significant speed-up

or at most a quadratic speed-up (e.g., Grover is optimal)

- ▶ NP-complete problems form a famous and important class of hard computational problems: satisfiability, Traveling Salesman Problem, protein folding, . . .

Conjectured: quantum computers can't efficiently solve them

# Conclusion

- ▶ Quantum mechanics is the **best physical theory we have**
- ▶ Fundamentally different from classical physics:  
superposition, interference, entanglement
- ▶ **Quantum algorithms** use these non-classical effects to solve some problems much faster
- ▶ We saw 4 important examples:
  1. Shor's factoring algorithm
  2. Grover's search algorithm
  3. Ambainis's collision-finding algorithm
  4. HHL algorithm for linear systems

**Much more left to be discovered. . .**

# Phase estimation

- ▶ Suppose we can apply  $U$  and are given one of its eigenvectors  $|v\rangle$  as a quantum state. Goal: learn eigenvalue  $e^{2\pi i\theta}$   
Suppose phase  $\theta = 0.\theta_1 \dots \theta_\ell$  has  $\ell$  bits of precision

- ▶ Remember QFT:  $|j\rangle \mapsto |\chi_j\rangle = \frac{1}{\sqrt{2^\ell}} \sum_{k=0}^{2^\ell-1} e^{\frac{2\pi ijk}{2^\ell}} |k\rangle$

- ▶ Phase estimation algorithm:

1. Start with  $|0^\ell\rangle|v\rangle$
2. Apply  $H^{\otimes \ell}$ :  $\frac{1}{\sqrt{2^\ell}} \sum_{k \in \{0,1\}^\ell} |k\rangle|v\rangle$
3. Conditioned on 1st register, apply  $U^k$  to 2nd register:

$$\frac{1}{\sqrt{2^\ell}} \sum_{k \in \{0,1\}^\ell} |k\rangle e^{2\pi i\theta k} |v\rangle = \frac{1}{\sqrt{2^\ell}} \sum_{k \in \{0,1\}^\ell} e^{2\pi i\theta k} |k\rangle |v\rangle$$

4. Inverse QFT on first register gives  $j = \theta 2^\ell = \theta_1 \dots \theta_\ell$

- ▶ With  $O(1/\varepsilon)$  applications of  $U$ :  $\varepsilon$ -error approximation of  $\theta$